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# Orthogonality Catastrophe

Martes Cuantico, April '17

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Orthogonality catastrophe:

$$\langle \text{GS}' | \text{GS} \rangle \rightarrow 0, \quad N \rightarrow \infty$$

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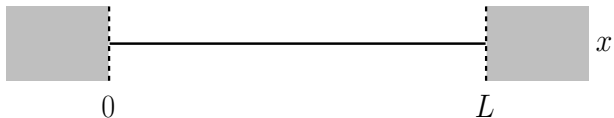
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$$\mathcal{A}_{\omega, \omega'} = \left[ \frac{4\omega/\omega'}{(1 + \omega/\omega')^2} \right]^{N/4} = e^{-\xi N}; \quad \xi > 0$$

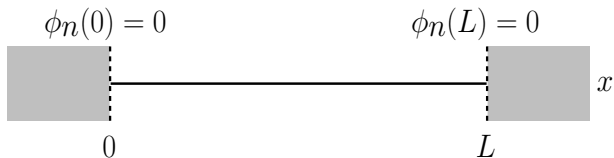
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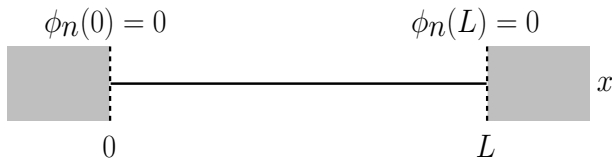


For each fermion

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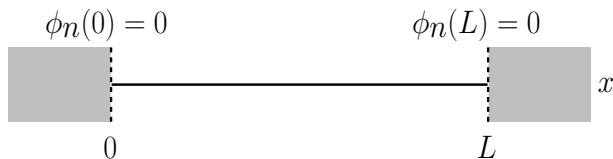
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$N$ -fermion ground state  $\Rightarrow$  Slater determinant

$$\Phi_N = \frac{1}{\sqrt{N!}} \det(\phi_n(x_j))_{n,j \leq N}$$

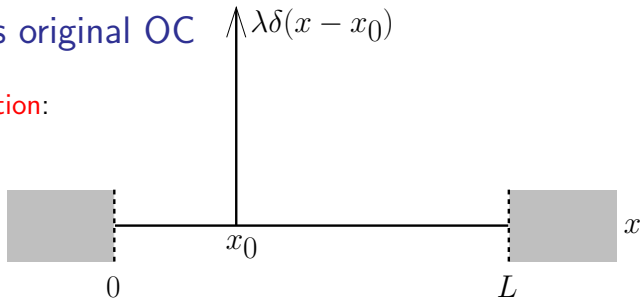


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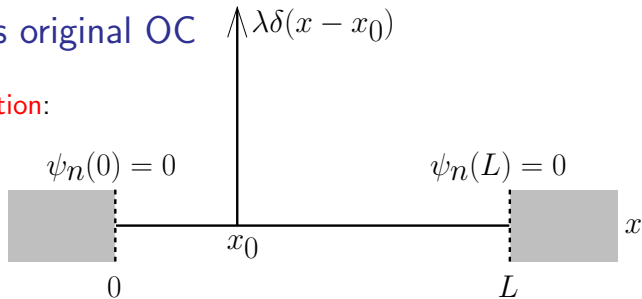
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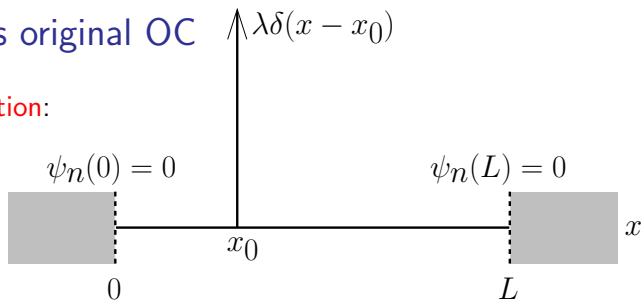


Now for a single fermion

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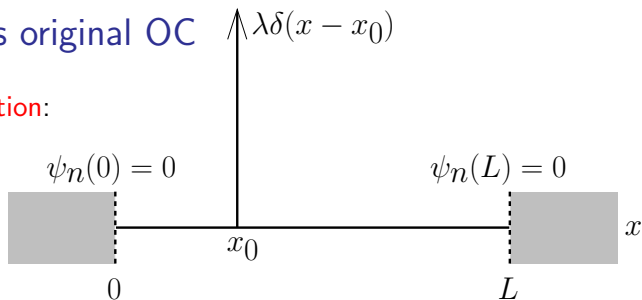


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Overlap of states before and after turn on  $\lambda\delta(x - x_0)$

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$$k'_n \approx k_n + \delta_n/L; \quad \delta_n = \lambda \frac{\sin^2(k_n x_0)}{k_n}$$

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If  $x_0 \ll L$

$$\mathcal{A}_{n,n'}^{(1)} \approx \frac{\pi n}{\pi(n + n') + \delta_{n'}} \frac{\sin \delta_{n'}}{\pi(n' - n) + \delta_{n'}}$$



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The  $N$ -particle overlap is bounded by the 1-particle overlaps

$$\mathcal{A} \leq \exp\left[-\sum_{\substack{n>N \\ n'<N}} |\mathcal{A}_{nn'}^{(1)}|^2\right]$$

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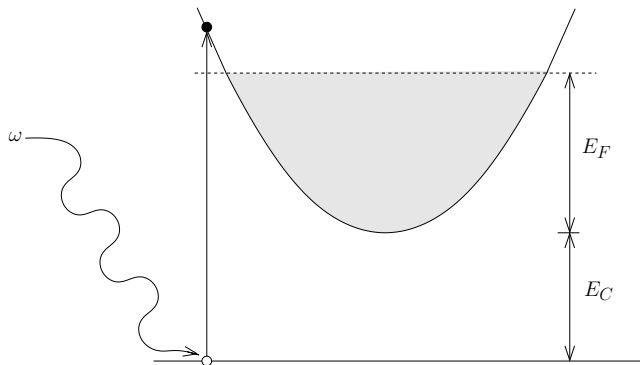
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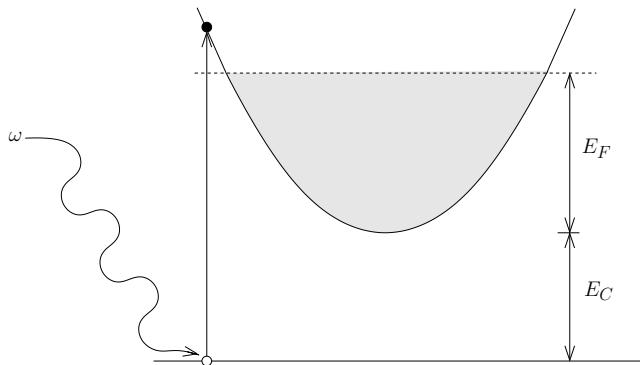
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Photon of frequency  $\omega \geq E_F + E_C$  excites a core  $e^-$  of a metal to the conduction band



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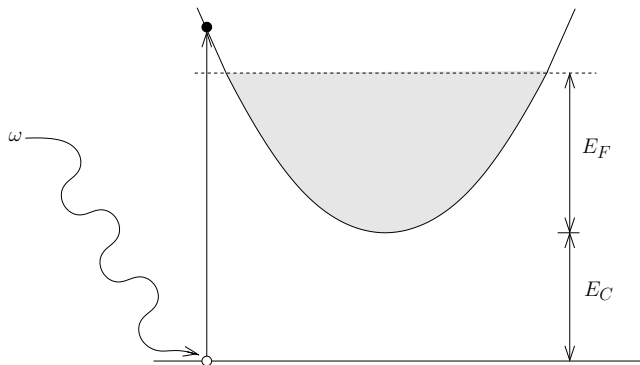
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OC between states of the Fermi sea before and after the absorption!

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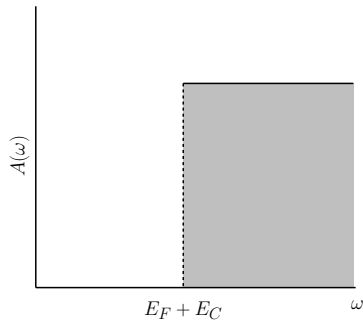
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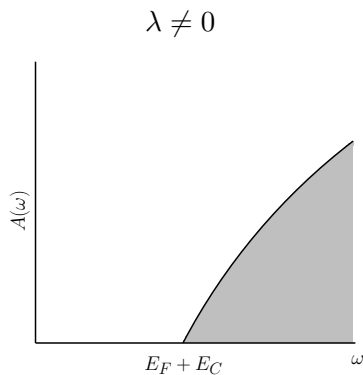
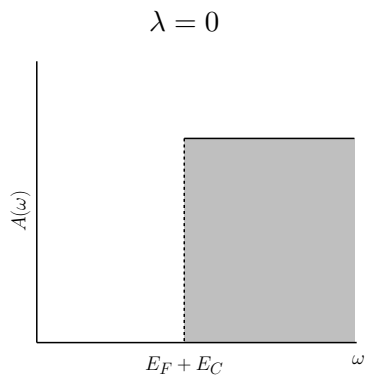
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$$\lambda = 0$$



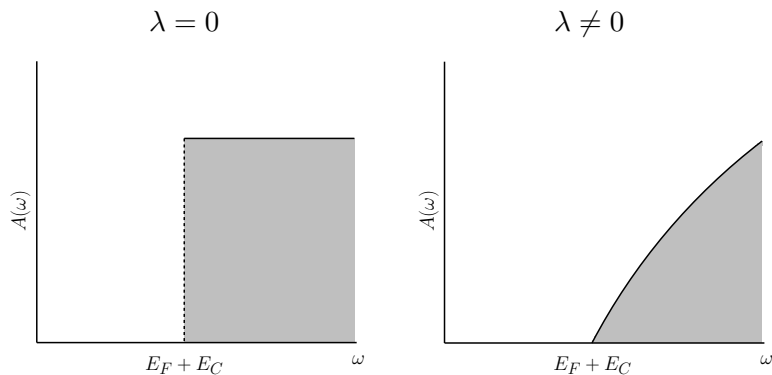
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OC tends to suppress the X-ray absorption

## Application II: Kondo effect

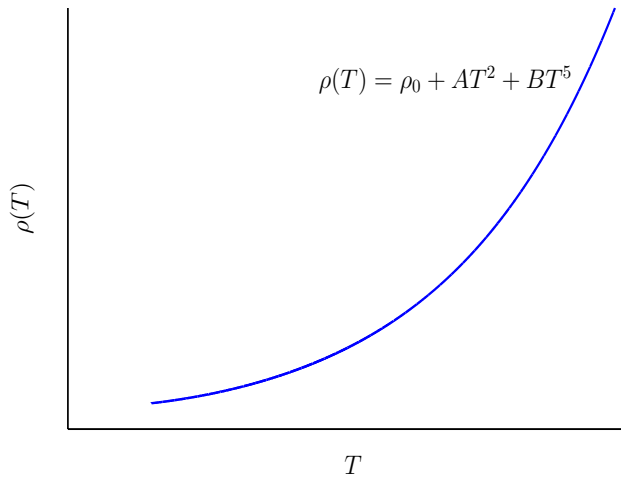
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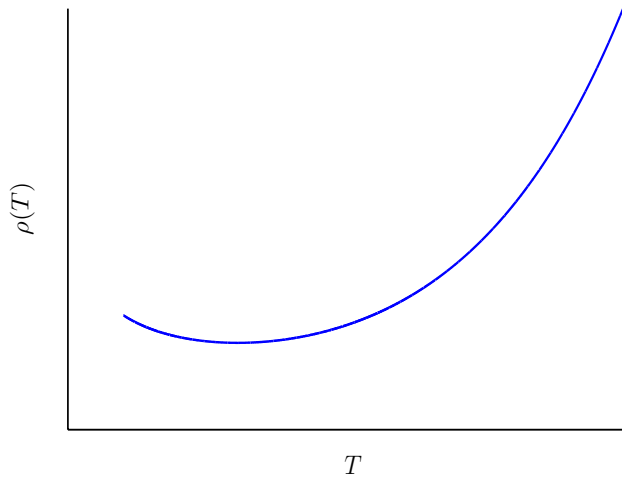
In general, electrical resistivity  $\rho$  of metals grows with temp.  $T$



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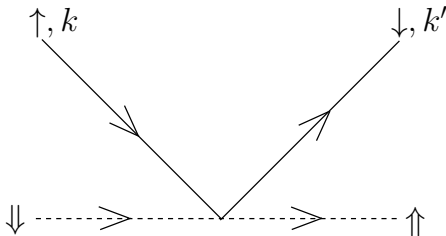
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In some cases, e.g. in gold,  $\rho(T)$  shows a minimum



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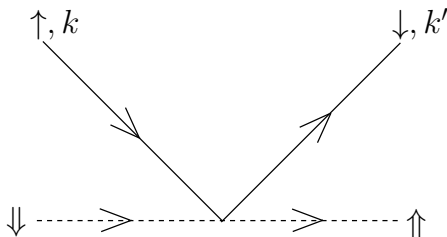
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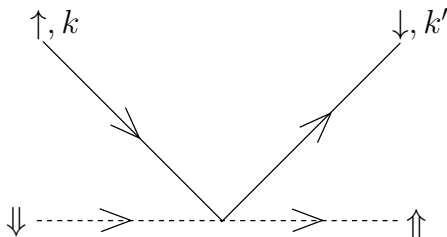


Perturbative computation  $\Rightarrow$  logarithmic contribution to  $\rho(T)$

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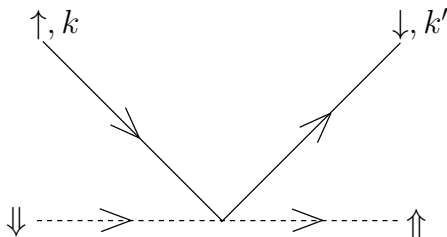
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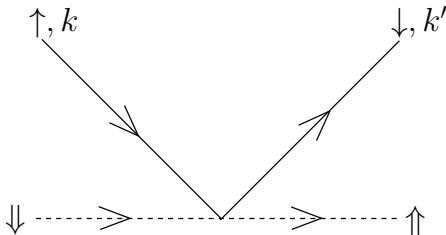
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Below Kondo temperature: perturbative methods are not valid

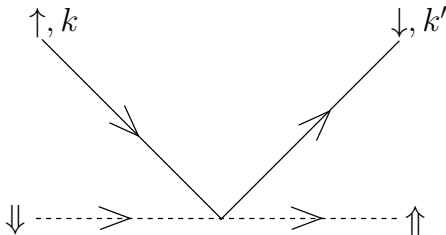
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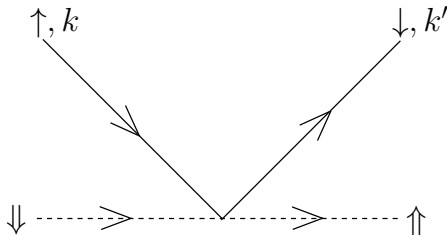
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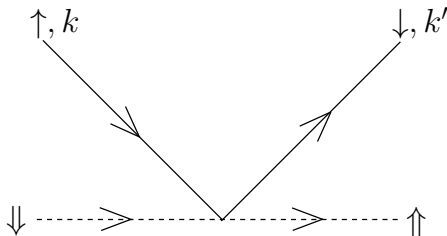


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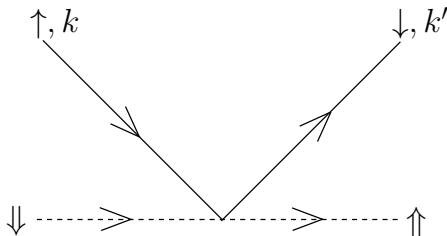


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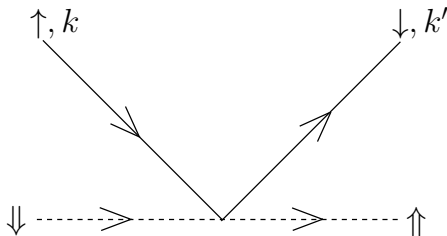
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**Kondo temperature:** scale at which spin-flips dominate over OC

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This system exhibits Haldane's exclusion statistics

## What is Haldane's exclusion statistics?

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$$E_N = \omega \sum_{j=1}^N n_j + \omega \frac{N}{2}$$

$n_j$ : energy level  $j$ -particle  $\Rightarrow n_j = 0, 1, 2, \dots$

$n_1 \leq n_2 \leq \dots \leq n_N \Rightarrow$  bosons



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If  $\mu \neq 0 \Rightarrow N$  interacting Harmonic oscillators with energy:

$$E_N = \omega \sum_{j=1}^N n_j + \omega \lambda \frac{N(N-1)}{2} + \omega \frac{N}{2}$$

$n_j$ : energy level  $j$ -particle  $\Rightarrow n_j = 0, 1, 2, \dots$

$n_1 \leq n_2 \leq \dots \leq n_N \Rightarrow$  (interacting) bosons

## What is Haldane's exclusion statistics?

$$H(\omega, \mu) = \frac{1}{2} \sum_{j=1}^N \left( -\partial_{x_j}^2 + \omega^2 x_j^2 \right) + \sum_{j=2}^N \sum_{k=1}^{j-1} \frac{\mu}{(x_j - x_k)^2}$$

However, defining

$$\tilde{n}_j = n_j + \lambda(j-1)$$

we have the spectrum of  $N$  free quasiparticles

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obeying the occupation rules

$$\tilde{n}_j \leq \tilde{n}_{j+1} - \lambda \Rightarrow \text{exclusion statistics}$$

# OC in the Sutherland model

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Take the ground state

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We have to solve a Selberg's integral...



## OC in the Sutherland model

Selberg's integral: higher-dimensional generalization of Euler  $B$  function

$$\int_0^1 \cdots \int_0^1 \prod_{i=1}^N x_i^{\alpha-1} (1-x_i)^{\beta-1} \prod_{1 \leq i, j \leq N} |x_i - x_j|^{2\gamma} dx_1 \cdots dx_N =$$
$$= \prod_{j=0}^{N-1} \frac{\Gamma(\alpha + j\gamma)\Gamma(\beta + j\gamma)\Gamma(1 + (j+1)\gamma)}{\Gamma(\alpha + \beta + (N+j-1)\gamma)\Gamma(1 + \gamma)}$$

$$\alpha, \beta, \gamma \in \mathbb{C}$$

$$\operatorname{Re}(\alpha), \operatorname{Re}(\beta) > 0$$

$$\operatorname{Re}(\gamma) > -\min\{1/N, \operatorname{Re}(\alpha)/(N-1), \operatorname{Re}(\beta)/(N-1)\}$$

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Exponential OC with  $N^2$

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Sutherland model: free quasiparticles obeying exclusion statistics

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- \* In localized **disordered systems**  
V Khemani *et al*, Nature Physics 11, 560 (2015)  
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¡Muchas gracias!